

GROUP CLASSIFICATION FOR A GENERAL NONLINEAR MODEL OF OPTION PRICING¹

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Abstract: We consider a family of equations with two free functional parameters containing the classical Black–Scholes model, Schönbucher–Wilmott model, Sircar–Papanicolaou equation for option pricing as partial cases. A five-dimensional group of equivalence transformations is calculated for that family. That group is applied to a search for specifications' parameters specifications corresponding to additional symmetries of the equation. Seven pairs of specifications are found.

Key words: Nonlinear partial differential equation, Group analysis, Group of equivalency transformations, Group classification, Nonlinear Black–Scholes equation, Pricing options, Dynamic hedging, Feedback effects of hedging.

Introduction

In the paper a nonlinear model

$$u_t + \frac{w(t, x)u_{xx}}{2(1 - xv(u_x)u_{xx})^2} + r(xu_x - u) = 0. \quad (0.1)$$

from the theory of financial markets is considered. In the case of $v \equiv 0$ it is generalized Black–Scholes equation [1], if, besides, $w(t, x) = \sigma^2 x^2$ (0.1) is the classical Black–Scholes model [2]. For arbitrary v and $w(t, x) = \sigma^2 x^2$ (0.1) is the Sircar–Papanicolaou nonlinear feedback pricing equation [1]. If v is arbitrary, $w(t, x) = \sigma^2 x^2$ and $r = 0$, it is the equilibrium pricing model or Schönbucher–Wilmott nonlinear feedback pricing model [3–6]. The last two models take into account a feedback effect of the presence of two types of traders. The programm traders are the portfolio insurers and the reference traders are the Black–Scholes uploaders.

The aim of the paper is to obtain a group classification [7] of equation (0.1) with free parameters v and w . The group of equivalence transformations [7, 8] of equation (0.1) will be found. By means of this group symmetries for the equation with all specifications will be calculated. Further these results will be applied to the theory of financial markets, particularly, they will allow to calculate various exact solutions of equation (0.1).

The groups of classical Black–Scholes model and their accordance to the groups of the heat equation were found in [9]. Research of symmetries of Schönbucher–Wilmott model and of some other nonlinear pricing models was made in [10–13].

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1. Group of the equivalence transformations

Let us find the continuous group of equivalence transformations of equation (0.1) for the applying to the search of specifications of the functions $v = v(u_x)$, $w = w(t, x)$ in the equation, that corresponds to additional symmetries for the symmetries of the kernel of principal Lie group for the equation. We rewrite equation (0.1) in the form

$$u_t + \frac{wu_{xx}}{2(1 - xvu_{xx})^2} + r(xu_x - u) = 0, \quad (1.1)$$

where v, w are the additional variables, depending on t, x, u, u_t and u_x . Generators of a continuous group of equivalence transformations will be searched in the form $Y = \tau \partial_t + \xi \partial_x + \eta \partial_u + \mu \partial_v + \nu \partial_w$, where the functions τ, ξ, η depend on t, x, u , and μ, ν depend on t, x, u, u_t, u_x, v, w . For brevity hereafter $\frac{\partial}{\partial t} \equiv \partial_t$ and similar notations are used. We add to (1.1) the equations

$$v_t = 0, \quad v_x = 0, \quad v_u = 0, \quad v_{u_t} = 0, \quad (1.2)$$

$$w_u = 0, \quad w_{u_t} = 0, \quad w_{u_x} = 0, \quad (1.3)$$

meaning that in the statement of the problem the function v depends only on u_x and the function w depends on t, x .

We consider the system of equations (1.1)–(1.3) as a manifold \mathfrak{N} in an expanded space of corresponding variables. Let us act on the left-hand side of system (1.1)–(1.3) by the extended operator

$$\tilde{Y} = Y + \varphi^t \partial_{u_t} + \varphi^{xx} \partial_{u_{xx}} + \mu^t \partial_{v_t} + \mu^x \partial_{v_x} + \mu^u \partial_{v_u} + \mu^{u_t} \partial_{v_{u_t}} + \nu^u \partial_{w_u} + \nu^{u_t} \partial_{w_{u_t}} + \nu^{u_x} \partial_{w_{u_x}},$$

we restrict a result of the action on \mathfrak{N} and we obtain the equations

$$\begin{aligned} \varphi^t + \frac{v w u_{xx}^2 \xi}{(1 - x v u_{xx})^3} + \frac{w(1 + x v u_{xx}) \varphi^{xx}}{2(1 - x v u_{xx})^3} + \frac{u_{xx} \nu}{2(1 - x v u_{xx})^2} + \frac{w x u_{xx}^2 \mu}{(1 - x v u_{xx})^3} + \\ + r(u_x \xi + x \varphi^x - \eta) \Big|_{\mathfrak{N}} = 0, \end{aligned} \quad (1.4)$$

$$\mu^t|_{\mathfrak{N}} = 0, \quad \mu^x|_{\mathfrak{N}} = 0, \quad \mu^u|_{\mathfrak{N}} = 0, \quad \mu^{u_t}|_{\mathfrak{N}} = 0, \quad (1.5)$$

$$\nu^u|_{\mathfrak{N}} = 0, \quad \nu^{u_t}|_{\mathfrak{N}} = 0, \quad \nu^{u_x}|_{\mathfrak{N}} = 0. \quad (1.6)$$

From (1.2) and (1.3) it follows that

$$\begin{aligned} \tilde{D}_t &= \partial_t + w_t \partial_w + w_{tt} \partial_{w_t} + w_{tx} \partial_{w_x} + \dots, \quad \tilde{D}_x = \partial_x + w_x \partial_w + w_{tx} \partial_{w_t} + w_{xx} \partial_{w_x} + \dots, \\ \tilde{D}_u &= \partial_u, \quad \tilde{D}_{u_t} = \partial_{u_t}, \quad \tilde{D}_{u_x} = \partial_{u_x} + v'(u_x) \partial_v + v''(u_x) \partial_{v'(u_x)} + \dots, \\ \mu^t &= \mu_t + w_t \mu_w - v'(u_x) \varphi_t^x = \mu_t + w_t \mu_w - v'(u_x) (\eta_{tx} + u_x \eta_{tu} - u_t \tau_{tx} - u_t u_x \tau_{tu} - u_x \xi_{tx} - u_x^2 \xi_{tu}), \\ \mu^x &= \mu_x + w_x \mu_w - v'(u_x) \varphi_x^x = \mu_x + w_x \mu_w - v'(u_x) (\eta_{xx} + u_x \eta_{xu} - u_t \tau_{xx} - u_t u_x \tau_{xu} - u_x \xi_{xx} - u_x^2 \xi_{xu}), \\ \mu^u &= \mu_u - v'(u_x) \varphi_u^x = \mu_u - v'(u_x) (\eta_{xu} + u_x \eta_{uu} - u_t \tau_{xu} - u_t u_x \tau_{uu} - u_x \xi_{xu} - u_x^2 \xi_{uu}), \\ \mu^{u_t} &= \mu_{u_t} - v'(u_x) \varphi_{u_t}^x = \mu_{u_t} + v'(u_x) \tau_x + u_x v'(u_x) \tau_u, \\ \nu^u &= \nu_u - w_t \tau_u - w_x \xi_u, \quad \nu^{u_t} = \nu_{u_t}, \quad \nu^{u_x} = \nu_{u_x} + v'(u_x) \nu_v. \end{aligned}$$

Therefore, equations (1.5) and (1.6) have the form

$$\mu_t + w_t \mu_w - v'(u_x) (\eta_{tx} + u_x \eta_{tu} - u_t \tau_{tx} - u_t u_x \tau_{tu} - u_x \xi_{tx} - u_x^2 \xi_{tu})|_{\mathfrak{N}} = 0, \quad (1.7)$$

$$\mu_x + w_x \mu_w - v'(u_x)(\eta_{xx} + u_x \eta_{xu} - u_t \tau_{xx} - u_t u_x \tau_{xu} - u_x \xi_{xx} - u_x^2 \xi_{xu})|_{\mathfrak{H}} = 0, \quad (1.8)$$

$$\mu_u - v'(u_x)(\eta_{xu} + u_x \eta_{uu} - u_t \tau_{xu} - u_t u_x \tau_{uu} - u_x \xi_{xu} - u_x^2 \xi_{uu})|_{\mathfrak{H}} = 0, \quad (1.9)$$

$$\mu_{u_t} + v'(u_x) \tau_x + u_x v'(u_x) \tau_u = 0, \quad (1.10)$$

$$\nu_u - w_t \tau_u - w_x \xi_u = 0, \quad \nu_{u_t} = 0, \quad \nu_{u_x} + v'(u_x) \nu_v = 0. \quad (1.11)$$

By equality (1.1) equations (1.7)–(1.9) can be rewritten in the form

$$\begin{aligned} \mu_t + w_t \mu_w - v'(u_x) (\eta_{tx} + u_x \eta_{tu} - u_x \xi_{tx} - u_x^2 \xi_{tu} + \\ + \frac{w u_{xx} (\tau_{tx} + u_x \tau_{tu})}{2(1 - x v u_{xx})^2} + (r x u_x - r u)(\tau_{tx} + u_x \tau_{tu})) = 0, \end{aligned} \quad (1.12)$$

$$\begin{aligned} \mu_x + w_x \mu_w - v'(u_x) (\eta_{xx} + u_x \eta_{xu} - u_x \xi_{xx} - u_x^2 \xi_{xu} + \\ + \frac{w u_{xx} (\tau_{xx} + u_x \tau_{xu})}{2(1 - x v u_{xx})^2} + (r x u_x - r u)(\tau_{xx} + u_x \tau_{xu})) = 0, \end{aligned} \quad (1.13)$$

$$\begin{aligned} \mu_u - v'(u_x) (\eta_{xu} + u_x \eta_{uu} - u_x \xi_{xu} - u_x^2 \xi_{uu} + \\ + \frac{w u_{xx} (\tau_{xu} + u_x \tau_{uu})}{2(1 - x v u_{xx})^2} + (r x u_x - r u)(\tau_{xu} + u_x \tau_{uu})) = 0. \end{aligned} \quad (1.14)$$

By means of the equality

$$\begin{aligned} \varphi^{xx} = \eta_{xx} + 2u_x \eta_{xu} + u_x^2 \eta_{uu} + u_{xx} \eta_u - u_t \tau_{xx} - 2u_t u_x \tau_{xu} - 2u_{tx} \tau_x - u_t u_x^2 \tau_{uu} - \\ - 2u_x u_{tx} \tau_u - u_t u_{xx} \tau_u - u_x \xi_{xx} - 2u_x^2 \xi_{xu} - 2u_{xx} \xi_x - u_x^3 \xi_{uu} - 3u_x u_{xx} \xi_u \end{aligned}$$

equation (1.4) is rewritten as

$$\begin{aligned} \eta_t + u_t \eta_u - u_t \tau_t - u_t^2 \tau_u - u_x \xi_t - u_t u_x \xi_u + \frac{1}{2(1 - x v u_{xx})^3} (2v w u_{xx}^2 \xi + 2x w u_{xx}^2 \mu + \\ + u_{xx} \nu - x u_{xx}^2 v \nu + w(1 + x v u_{xx})(\eta_{xx} + 2u_x \eta_{xu} + u_x^2 \eta_{uu} + u_{xx} \eta_u - \\ - u_t \tau_{xx} - 2u_t u_x \tau_{xu} - 2u_{tx} \tau_x - u_t u_x^2 \tau_{uu} - 2u_x u_{tx} \tau_u - \\ - u_t u_{xx} \tau_u - u_x \xi_{xx} - 2u_x^2 \xi_{xu} - 2u_{xx} \xi_x - u_x^3 \xi_{uu} - 3u_x u_{xx} \xi_u)) + \\ + r u_x \xi + r x (\eta_x + u_x \eta_u - u_t \tau_x - u_t u_x \tau_u - u_x \xi_x - u_x^2 \xi_u) - r \eta|_{\mathfrak{H}} = \\ = \eta_t + \frac{w u_{xx} (\tau_t - \eta_u)}{2(1 - x v u_{xx})^2} + (r x u_x - r u)(\tau_t - \eta_u) - \frac{w^2 u_{xx}^2 \tau_u}{4(1 - x v u_{xx})^4} - \\ - (r x u_x - r u)^2 \tau_u + \frac{w u_{xx} (r x u_x - r u) \tau_u}{(1 - x v u_{xx})^2} - u_x \xi_t + \\ + \frac{w u_x u_{xx} \xi_u}{2(1 - x v u_{xx})^2} + (r x u_x - r u) u_x \xi_u + \frac{1}{2(1 - x v u_{xx})^3} (2v w u_{xx}^2 \xi + 2x w u_{xx}^2 \mu + \\ + u_{xx} \nu - x u_{xx}^2 v \nu + w(1 + x v u_{xx})(\eta_{xx} + 2u_x \eta_{xu} + u_x^2 \eta_{uu} + u_{xx} \eta_u + \\ + \frac{w u_{xx} \tau_{xx}}{2(1 - x v u_{xx})^2} + (r x u_x - r u) \tau_{xx} + \frac{w u_x u_{xx} \tau_{xu}}{(1 - x v u_{xx})^2} + 2(r x u_x - r u) u_x \tau_{xu} - \\ - 2u_{tx} \tau_x + \frac{w u_x^2 u_{xx} \tau_{uu}}{2(1 - x v u_{xx})^2} + (r x u_x - r u) u_x^2 \tau_{uu} - 2u_x u_{tx} \tau_u + \frac{w u_{xx}^2 \tau_u}{2(1 - x v u_{xx})^2} + \\ + (r x u_x - r u) u_{xx} \tau_u - u_x \xi_{xx} - 2u_x^2 \xi_{xu} - 2u_{xx} \xi_x - u_x^3 \xi_{uu} - 3u_x u_{xx} \xi_u)) + \\ + r u_x \xi + r x (\eta_x + u_x \eta_u - u_x \xi_x - u_x^2 \xi_u) - r \eta + \\ + \frac{r x w u_{xx} (\tau_x + u_x \tau_u)}{2(1 - x v u_{xx})^2} + r x (r x u_x - r u)(\tau_x + u_x \tau_u) = 0. \end{aligned} \quad (1.15)$$

We differentiate the last equations with respect to u_{tx} and obtain $w(1 + xvu_{xx})(\tau_x + u_x\tau_u) = 0$, consequently, $\tau = \tau(t)$, if $w \neq 0$. Therefore, equations (1.11)–(1.15) have the form

$$\mu_{u_t} = 0, \quad \nu_u - w_x\xi_u = 0, \quad \nu_{u_t} = 0, \quad \nu_{u_x} + v'(u_x)\nu_v = 0, \quad (1.16)$$

$$\mu_t + w_t\mu_w - v'(u_x)(\eta_{tx} + u_x\eta_{tu} - u_x\xi_{tx} - u_x^2\xi_{tu}) = 0, \quad (1.17)$$

$$\mu_x + w_x\mu_w - v'(u_x)(\eta_{xx} + u_x\eta_{xu} - u_x\xi_{xx} - u_x^2\xi_{xu}) = 0, \quad (1.18)$$

$$\mu_u - v'(u_x)(\eta_{xu} + u_x\eta_{uu} - u_x\xi_{xu} - u_x^2\xi_{uu}) = 0, \quad (1.19)$$

$$\begin{aligned} \eta_t + \frac{wu_{xx}(\tau'(t) - \eta_u)}{2(1 - xvu_{xx})^2} + (rxu_x - ru)(\tau'(t) - \eta_u) - u_x\xi_t + \\ + \frac{wu_xu_{xx}\xi_u}{2(1 - xvu_{xx})^2} + (rxu_x - ru)u_x\xi_u + \frac{1}{2(1 - xvu_{xx})^3} (2vwu_{xx}^2\xi + 2xwu_{xx}^2\mu + \\ + u_{xx}\nu - xu_{xx}^2v\nu + w(1 + xvu_{xx})(\eta_{xx} + 2u_x\eta_{xu} + u_x^2\eta_{uu} + u_{xx}\eta_u - u_x\xi_{xx} - \\ - 2u_x^2\xi_{xu} - 2u_{xx}\xi_x - u_x^3\xi_{uu} - 3u_xu_{xx}\xi_u)) + ru_x\xi + \\ + rx(\eta_x + u_x\eta_u - u_x\xi_x - u_x^2\xi_u) - r\eta = 0. \end{aligned}$$

We multiply by $2(1 - xvu_{xx})^3$ the last equation, then

$$\begin{aligned} 2(1 - xvu_{xx})^3(\eta_t + (rxu_x - ru)(\tau'(t) - \eta_u) + u_x\xi_u) + (1 - xvu_{xx})wu_{xx}(\tau'(t) - \eta_u) - \\ - 2(1 - xvu_{xx})^3u_x\xi_t + (1 - xvu_{xx})wu_xu_{xx}\xi_u + 2vwu_{xx}^2\xi + 2xwu_{xx}^2\mu + \\ + u_{xx}\nu - xu_{xx}^2v\nu + w(1 + xvu_{xx})(\eta_{xx} + 2u_x\eta_{xu} + u_x^2\eta_{uu} + u_{xx}\eta_u - u_x\xi_{xx} - \\ - 2u_x^2\xi_{xu} - 2u_{xx}\xi_x - u_x^3\xi_{uu} - 3u_xu_{xx}\xi_u) + \\ + 2(1 - xvu_{xx})^3(ru_x\xi + rx(\eta_x + u_x\eta_u - u_x\xi_x - u_x^2\xi_u) - r\eta) = 0. \end{aligned} \quad (1.20)$$

Equation (1.20) for the case $v \neq 0$ has at u_{xx}^3 multiplier

$$\eta_t + rxu_x\tau'(t) - ru(\tau'(t) - \eta_u + u_x\xi_u) - u_x\xi_t + ru_x\xi + rx(\eta_x - u_x\xi_x) - r\eta,$$

after its splitting with respect to u_x , we obtain two equations

$$\eta_t + rx\eta_x + ru\eta_u - r\eta - ru\tau'(t) = 0, \quad (1.21)$$

$$rx\tau'(t) - \xi_t - rx\xi_x - ru\xi_u + r\xi = 0. \quad (1.22)$$

After the splitting with respect to u_x of the multiplier at u_{xx} in zero degree it follows that

$$\xi = A(t, x)u + B(t, x),$$

$$\eta = A_x(t, x)u^2 + C(t, x)u + D(t, x)$$

and by (1.21), (1.22)

$$\begin{aligned} 2\eta_t - 2ru\tau' + 2ru\eta_u + 2rx\eta_x - 2r\eta + w\eta_{xx} = w\eta_{xx} = 0, \\ 2rx\tau' - 2ru\xi_u - 2\xi_t + 2r\xi - 2rx\xi_x - 2w\xi_{xx} + 4w\eta_{xu} = -w\xi_{xx} + 2w\eta_{xu} = 0. \end{aligned} \quad (1.23)$$

The last equality implies that

$$A_{xx} = 0, \quad A(t, x) = A_1(t)x + A_0(t), \quad C(t, x) = \frac{1}{2}B_x(t, x) + E(t),$$

$$\xi = A_1(t)xu + A_0(t)u + B(t, x), \quad \eta = A_1(t)u^2 + \frac{1}{2}B_x(t, x)u + E(t)u + D(t, x).$$

Then from (1.23) it follows that

$$\begin{aligned} B_{xxx} &= 0, \quad D_{xx} = 0, \quad \xi = A_1(t)xu + A_0(t)u + B_2(t)x^2 + B_1(t)x + B_0(t), \\ \eta &= A_1(t)u^2 + B_2(t)xu + \frac{1}{2}B_1(t)u + E(t)u + D_1(t)x + D_0(t). \end{aligned}$$

Now the equality (1.22) implies that $A_1(t) = Fe^{-rt}$, $A_0(t)$ is a constant,

$$\begin{aligned} B_2(t) &= Ge^{-rt}, \quad B_1(t) = r\tau(t) + H, \quad B_0(t) = Je^{rt}, \\ \xi &= Fe^{-rt}xu + A_0u + Ge^{-rt}x^2 + r\tau(t)x + Hx + Je^{rt}, \\ \eta &= Fe^{-rt}u^2 + Ge^{-rt}xu + \frac{1}{2}(r\tau(t) + H)u + E(t)u + D_1(t)x + D_0(t). \end{aligned}$$

By (1.21) D_1 is a constant,

$$\begin{aligned} D_0(t) &= Ke^{rt}, \quad E(t) = \frac{1}{2}r\tau(t) + P, \\ \eta &= Fe^{-rt}u^2 + Ge^{-rt}xu + r\tau(t)u + Pu + D_1x + Ke^{rt}. \end{aligned}$$

From (1.16) it follows that $\nu = w_x(Fe^{-rt}x + A_0)u + S(t, x, u_x, v, w)$.

The coefficient at u_{xx} in equation (1.20) is equated to zero and we obtain the equation

$$\begin{aligned} &-6xv\eta_t - 6xv(rxu_x\tau'(t) - ru\tau'(t) + ru\eta_u - ruu_x\xi_u) + w\tau'(t) + \\ &+ 6xvu_x\xi_t + wu_x\xi_u + \nu - 2w\xi_x - 3wu_x\xi_u + 2xvwu_x\eta_{xu} + xvwu_x^2\eta_{uu} - \\ &- xvwu_x\xi_{xx} - 2xvwu_x^2\xi_{xu} - 6xv(ru_x\xi + rx\eta_x - rxu_x\xi_x - r\eta) = 0. \end{aligned}$$

Let us substitute in it the expressions for ξ , η , ν that were found before, and splitting with respect to the variable u leads to the equations

$$-2Fe^{-rt}w + w_x(Fe^{-rt}x + A_0) = 0, \quad (1.24)$$

$$S = 4Ge^{-rt}xw - w\tau' + 2rw\tau + 2Hw + 2Fe^{-rt}xu_xw + 2A_0u_xw.$$

The last of them implies the equalities $\nu_v = S_v = 0$, consequently, by (1.16) we obtain

$$\nu_{u_x} = S_{u_x} = 2Fe^{-rt}xw + 2A_0xw = 0, \quad A_0 = F = 0.$$

Thus,

$$\begin{aligned} \xi &= Ge^{-rt}x^2 + r\tau(t)x + Hx + Je^{rt}, \\ \eta &= Ge^{-rt}xu + r\tau(t)u + Pu + D_1x + Ke^{rt}, \\ \nu &= 4Ge^{-rt}xw - w\tau' + 2rw\tau + 2Hw. \end{aligned}$$

Analogous calculations are made with the coefficient at u_{xx}^2 in equation (1.20), we obtain the equation

$$\begin{aligned} &6x^2v^2(\eta_t + rxu_x\tau' - ru\tau' + ru\eta_u) - xvw(\tau' - \eta_u) - 6x^2v^2u_x\xi_t + 2vw\xi + 2xw\mu - xvv + \\ &+ xvw(\eta_u - 2\xi_x) + 6x^2v^2(ru_x\xi + rx\eta_x - rxu_x\xi_x - r\eta) = 0, \end{aligned}$$

that implies the equality $\mu = v(H - P - \frac{J}{x}e^{rt} + 2Ge^{-rt}x)$. Therefore $\mu_u = \mu_w = 0$, and for the case $v' \neq 0$ obtain $G = 0$ from equation (1.19). Then equation (1.18) implies that $\mu_x = 0$, hence $J = 0$. From equation (1.17) it follows that $\mu_t = 0$, it corresponds to the resulting formula $\mu = (H - P)v$. Thus, $\tau(t)$ is an arbitrary function,

$$\begin{aligned} \xi &= Hx + r\tau(t)x, \quad \eta = Ke^{rt} + D_1x + Pu + r\tau(t)u, \\ \mu &= (H - P)v, \quad \nu = 2Hw + (2r\tau(t) - \tau'(t))w. \end{aligned}$$

Let us formulate the result in the form of theorem.

Theorem 1. *The Lie algebra of infinitesimal generators of the equivalency transformations groups for equation (0.1), is generated by operators*

$$Y_1 = x\partial_u, \quad Y_2 = e^{rt}\partial_u, \quad Y_3 = x\partial_x + u\partial_u + 2w\partial_w, \quad Y_4 = x\partial_x + v\partial_v + 2w\partial_w, \\ Y_5 = \tau(t)\partial_t + r\tau(t)x\partial_x + r\tau(t)u\partial_u + (2r\tau(t) - \tau'(t))w\partial_w,$$

when v', w are identically unequal to zero.

Remark 1. It is easy to check that the infinitely-dimensional part of the Lie algebra from Theorem 1 consists of operators of the form Y_5 only.

The extensions of the operators Y_k , $k = 1, 2, 3, 4, 5$, are

$$\tilde{Y}_1 = x\partial_u + \partial_{u_x}, \quad \tilde{Y}_2 = e^{rt}\partial_u, \quad \tilde{Y}_3 = x\partial_x + u\partial_u + 2w\partial_w, \\ \tilde{Y}_4 = x\partial_x + v\partial_v + 2w\partial_w - u_x\partial_{u_x}, \quad \tilde{Y}_5 = \tau\partial_t + r\tau x\partial_x + r\tau u\partial_u + (2r\tau - \tau')w\partial_w. \quad (1.25)$$

Therefore, the kernel of the principal Lie algebras for equation (0.1) is one-dimensional with the basis Y_2 , because the corresponding group only doesn't transform the additional variables v, w and their arguments t, x, u_x .

Corollary 1. *The kernel of the principal Lie algebras for equation (0.1) is spanned by the operator $X_1 = e^{rt}\partial_u$ when v', w are identically unequal to zero.*

2. Group classification

Consider Lie algebra of projections of operators (1.25) on the subspace of the variables t, x, u_x, v, w , i. e. the algebra generated by

$$Z_1 = \partial_{u_x}, \quad Z_2 = v\partial_v - u_x\partial_{u_x}, \\ Z_3 = x\partial_x + 2w\partial_w, \quad Z_4 = \tau\partial_t + r\tau x\partial_x + (2r\tau - \tau')w\partial_w. \quad (2.1)$$

It is the direct sum of subalgebras $\langle Z_1, Z_2 \rangle$ and $\langle Z_3, Z_4 \rangle$ that corresponds to two different functions v and w and their different arguments. Therefore, the subalgebras can be considered separately.

Nonzero structure constants of $\langle Z_1, Z_2 \rangle$ are $c_{12}^1 = -1$, $c_{21}^1 = 1$. Therefore, the inner automorphisms are $E_1 : \bar{e}^1 = e^1 - e^2 a_1$, $E_2 : \bar{e}^1 = e^1 e^{a_2}$. Here e^i , $i = 1, 2$ are the coefficients at Z_i respectively in the basis decomposition of Z . If $e^2 \neq 0$, then $e^1 = 0$ by the acting of E_1 . Therefore the optimal system of one-dimensional subalgebras consists of subalgebras with bases Z_1 and Z_2 .

In the subalgebra $\langle Z_3, Z_4 \rangle$ there are no nontrivial inner automorphisms, consequently, the optimal system of one-dimensional subalgebras has a form $\Theta_1 = \{\langle Z_2 \rangle, \langle bZ_2 + Z_4 \rangle, b \in \mathbb{R}\}$.

For operators Z from optimal systems we calculate the expressions

$$Z(V(u_x) - v)|_{v=V} = 0, \quad Z(W(t, x) - w)|_{w=W} = 0.$$

Note, that if Z contains Z_1 with a nonzero coefficient and doesn't contain Z_2 , then $v' = 0$. Such case doesn't correspond to the conditions of Theorem 1. If an operator Z has nonzero coefficients at Z_1 and at Z_3 , then by E_1 the coefficient at Z_1 can be equated to zero for equivalent operator to Z . Therefore, the operator Z_1 can be excluded from further considerations.

We have

$$Z_2(V(u_x) - v)|_{w=W} = -V - u_x V' = 0, \quad V = \beta/u_x$$

for arbitrary $\beta \in \mathbb{R}$. Further,

$$Z_3(W(t, x) - w)|_{w=W} = xW_x - 2W = 0, \quad W = D(t)x^2$$

for arbitrary function $D(t)$. Finally,

$$(bZ_3 + Z_4(W(t, x) - w)|_{w=W} = \tau(t)W_t + (r\tau(t) + b)xW_x - (2r\tau(t) - \tau'(t) + 2b)W = 0,$$

$$W = \frac{e^{2rt+2b \int \frac{dt}{\tau(t)}}}{\tau(t)} \varphi(xe^{-rt-b \int \frac{dt}{\tau(t)}})$$

for arbitrary functions $\varphi \neq 0$, $\tau \neq 0$.

Optimal system of two-dimensional subalgebras consists of $\langle Z_2, Z_3 \rangle$, $\langle Z_2, bZ_3 + Z_4 \rangle$, $\langle Z_3, Z_4 \rangle$. In the first two cases we have the simultaneous specifications for v and w that are already known. In the last one specification we have the form $W = \gamma x^2 / \tau(t)$.

For the Lie algebra $\langle Z_2, Z_3, Z_4 \rangle$ the specifications are $V = \beta / u_x$, $W = \gamma x^2 / \tau(t)$.

For every basis operator from the optimal systems calculate the projection of the corresponding generator of the group of equivalency transformations on the space of the variables t, x, u . Then Z_2 corresponds to $\text{pr}_{(t,x,u)}(Y_4 - Y_3) = -u\partial_u$, for the operator Z_3 it will be $\text{pr}_{(t,x,u)}Y_3 = x\partial_x + u\partial_u$, and $\text{pr}_{(t,x,u)}(bY_3 + Y_5) = \tau(t)\partial_t + (r\tau(t) + b)x\partial_x + (r\tau(t)u + b)\partial_u$ corresponds to $bZ_3 + Z_4$. It implies the next theorem.

Theorem 2. *Let v' , w be identically unequal to zero i. Then next assertions are true.*

1. *The principal Lie algebra of the equation*

$$u_t + \frac{w(t, x)u_{xx}}{2 \left(1 - \beta x \frac{u_{xx}}{u_x}\right)^2} + r(xu_x - u) = 0, \quad \beta \neq 0,$$

is generated by the operators $X_1 = e^{rt}\partial_u$, $X_2 = u\partial_u$.

2. *The principal Lie algebra of the equation*

$$u_t + \frac{T'(t)e^{2rt+2bT(t)}\varphi(xe^{-rt-bT(t)})u_{xx}}{2(1 - xv(u_x)u_{xx})^2} + r(xu_x - u) = 0, \quad T'(t) \neq 0, \quad \varphi(z) \neq 0,$$

is generated by the operators $X_1 = e^{rt}\partial_u$, $X_2 = \frac{1}{T'(t)}\partial_t + \left(\frac{r}{T'(t)} + b\right)x\partial_x + \left(\frac{r}{T'(t)} + b\right)u\partial_u$.

3. *The principal Lie algebra of the equation*

$$u_t + \frac{T'(t)e^{2rt+2bT(t)}\varphi(xe^{-rt-bT(t)})u_{xx}}{2 \left(1 - \beta x \frac{u_{xx}}{u_x}\right)^2} + r(xu_x - u) = 0, \quad T'(t) \neq 0, \quad \varphi(z) \neq 0, \quad \beta \neq 0,$$

is generated by the operators

$$X_1 = e^{rt}\partial_u, \quad X_2 = u\partial_u, \quad X_3 = \frac{1}{T'(t)}\partial_t + \left(\frac{r}{T'(t)} + b\right)x\partial_x + \left(\frac{r}{T'(t)} + b\right)u\partial_u.$$

4. *The principal Lie algebra of the equation*

$$u_t + \frac{D(t)x^2u_{xx}}{2(1 - xv(u_x)u_{xx})^2} + r(xu_x - u) = 0, \quad D(t) \neq 0,$$

is generated by the operators

$$X_1 = e^{rt}\partial_u, \quad X_2 = x\partial_x + u\partial_u, \quad X_3 = \frac{1}{D(t)}\partial_t + \frac{r}{D(t)}x\partial_x + \frac{r}{D(t)}u\partial_u.$$

5. The principal Lie algebra of the equation

$$u_t + \frac{D(t)x^2u_{xx}}{2\left(1 - \beta x \frac{u_{xx}}{u_x}\right)^2} + r(xu_x - u) = 0, \quad D(t) \neq 0,$$

is generated by the operators

$$X_1 = e^{rt}\partial_u, \quad X_2 = x\partial_x, \quad X_3 = u\partial_u, \quad X_4 = \frac{1}{D(t)}\partial_t + \frac{r}{D(t)}x\partial_x + \frac{r}{D(t)}u\partial_u.$$

Remark 2. Theorem 1 and Theorem 2 are valid for the case $r = 0$.

3. Conclusion

Further Theorem 2 will be applied to the search of exact solutions of the option pricing nonlinear models. Specification $W(t, x) = \sigma^2 x^2$ as partial case of $D(t)x^2$ corresponds to the Scönbucher—Wilmott model, if $r = 0$, and to Círcar—Papanicolaou model for $r \neq 0$.

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